

## THE KEPLER PROBLEM

Due date: Jan 31, 2024 — 9am

A two-body problem in a central potential can be described analytically. The orbital equation cannot be written explicitly as a function of time ( $dr/dt$ ), but it can be written explicitly as a function of angle ( $dr/dv$ ). In practice, however, we frequently need to compute the location of a moving object as a function of time. In such cases we resort to two possible approaches: solve the problem iteratively, with anomalies, or solve the differential equation numerically. We will apply both methods to the motion of Halley's comet and to study the long-term (numerical) stability of its orbit.

1. **Iterative solution.** Halley's comet might be the most famous comet of all time. Its perihelion distance is 0.587 au, eccentricity is 0.967 and the orbital period is 76 years. Using the Newton-Raphson method, compute the position of Halley's comet around the Sun as a function of time. Plot the orbit.
2. **Runge-Kuta approximation.** Compute the coefficients  $\alpha_1, \alpha_2, \beta_1, \beta_2, \beta_3, \gamma_1, \gamma_2$  and  $\gamma_3$  for the third order of Runge-Kutta expansion. Typeset the derivation, mostly for L<sup>A</sup>T<sub>E</sub>X practice.
3. **Differential solution.** Using an off-the-shelf 4<sup>th</sup> order Runge-Kutta method, compute the orbit of Halley's comet and compare it with the orbit obtained by the iterative solution. Plot the accumulated error as a function of time (or the number of orbits).
4. **Stability.** Study the stability of Halley's comet using several differential integrators. Be sure to explicitly reference the used code or algorithms.

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A quick recap: the equation of an elliptical orbit is:

$$r(v) = \frac{a(1 - e^2)}{1 + e \cos v}, \quad (1)$$

where  $a$  is the semi-major axis,  $e$  is orbital eccentricity and  $v$  is true anomaly. The true anomaly is measured from perihelion and can be related to time via the eccentric and mean anomalies:

$$\tan \frac{v}{2} = \sqrt{\frac{1+e}{1-e}} \tan \frac{E}{2}, \quad M = E - e \sin E = \frac{2\pi(t - \tau)}{P}, \quad (2)$$

where  $E$  is the eccentric anomaly,  $M$  is the mean anomaly,  $\tau$  is the reference time at perihelion and  $P$  is the orbital period. Equation  $M = E - e \sin E$  cannot be solved analytically for  $E$ , thus it needs to be solved iteratively, using the Newton-Raphson method. This is also the reason why  $r(t)$  cannot be written out explicitly.